MTH 654-2024 Homework C2. Date: January 12, 2025 Prof. M. Peszynska

1 Introduction.

In this project, we consider optimal control problems under PDE constraints. In particular, we focus on boundary and distributed control problems constrained by the Poisson equation. The work was guided by [1] and [2].

2 An uncoupled boundary control problem.

The first part of the project involved solving the following problem:

Problem 2.1 – 1D uncoupled boundary control problem

Suppose we have a 1-dimensional rod represented by the interval [0, 1]. We will control the heat on the left side of the rod by $u \in \mathbb{R}$ and fix the rod's temperature at the right endpoint to be 0. We would like the steady-state temperature v of each point x of the rod to be as close to 1 as possible. From physics, we know v must satisfy $-\Delta v = 0$ in $\Omega = (0, 1)$, which motivates solving the following constrained optimization problem:

$$\underset{u,v}{\text{minimize }} J(v,u) := \int_{\Omega} (v(x) - 1)^2 \, dx + \frac{\lambda}{2} u^2 \tag{1a}$$

subject to
$$-\Delta v = 0$$
 on Ω , (1b)

 $v(0) = u, \ v(1) = 0 \tag{1c}$

$$u \in \mathbb{R}.$$
 (1d)

Let's introduce some typical terminology for constrained optimization problems of this form: u is the *control*, v is the *state*, the function J being minimized in (1a) is the *objective*, $0 \leq \lambda \ll 1$ is a fixed *regularization parameter*, (1b)-(1d) are called the *constraints*, and \mathbb{R} in (1d) defines the set of *admissible controls*.

We will solve this problem two ways. First, we solve the problem analytically using the fact that the control u comes from a finite-dimensional vector space. Then, we solve the problem numerically using a method that is applied to higher-dimensional problems, where the control is generally taken to be a function $u \in L^2(\partial\Omega)$.

2.1 Analytic solution.

As posed in (1), the objective is a function of two variables, the control u and the state v. From the theory of differential equations, though, we know that each $u \in \mathbb{R}$ gives rise to a unique solution v that satisfies the constraints of the problem. So, we begin by defining the solution operator, $S : \mathbb{R} \to L^2(\Omega)$, such that $u \mapsto v$. This means we may re-frame the problem as

$$\underset{u}{\text{minimize } j(u)} := \int_{\Omega} (S(u)(x) - 1)^2 \, dx + \frac{\lambda}{2} u^2 \tag{2a}$$

subject to
$$u \in \mathbb{R}$$
. (2b)

To solve this problem, we compute a formula for the solution operator. From the theory of differential equations, if v solves (1b), then v(x) = ax + b for some $a, b \in \mathbb{R}$. Inputting the boundary conditions in (1c) therefore yields the formula v(x) = u(1-x), so that S(u)(x) = u(1-x) is the solution operator we wanted. Thus the modified objective becomes

$$j(u) := \int_{\Omega} (S(u)(x) - 1)^2 dx + \frac{\lambda}{2}u^2$$
$$= \int_0^1 (u(1-x) - 1)^2 dx + \frac{\lambda}{2}u^2$$
$$= 1 - u + \frac{1}{3}u^2 + \frac{\lambda}{2}u^2.$$

We therefore solve j'(u) = 0 and take $\lambda \to 0$ to get u = 3/2.

2.2 Numerical solution using scipy.optimize.

To solve the problem numerically, we follow an approach that applies even when u comes from an infinite-dimensional function space as is the case for higher dimensions. We begin by discretizing everything in the problem: given $M \in \mathbb{N}$, we partition Ω into a grid with boundaries $0 = x_{1/2} < x_{3/2} < \cdots < x_{M+1/2} = 1$, centers $x_k := 1/2 \cdot (x_{k-1/2} + x_{k+1/2})$, and widths $h_k := x_{k+1/2} - x_{k-1/2}$. Then define $\omega_k := [x_{k-1/2}, x_{k+1/2}]$, and approximate v by $v_h = \sum_{k=1}^{M} v_k \mathbb{1}_{\omega_k}$. Under this discretization, we reformulate the objective and constraints to get

minimize
$$J_h(v_h, u) := \frac{1}{2} \int_{\Omega} (v_h(x) - 1)^2 dx + \frac{\lambda}{2} u^2$$
 (3a)

subject to
$$\frac{1}{h_k^2}(-v_{k-1}+2v_k-v_{k+1})=0$$
 for $k=2,\ldots,M-1,$ (3b)

$$v_1 = u, \ v_m = 0 \tag{3c}$$

$$u \in \mathbb{R}.$$
 (1d)

where, in general, the differential equation constraint in (3b) comes from a discrete version of the Laplacian operator Δ . To solve this problem using **scipy.optimize**, we simplify the objective to get

$$J_h(v_h, u) = \frac{1}{2} \sum_{k=1}^{M} h_k \cdot (v_k - 1)^2 + \frac{\lambda}{2} v_1^2$$

and we rewrite the constraints in matrix-vector form as

$$\begin{bmatrix} 0 & 0 & 0 & 0 & \cdots & 0 \\ -1 & 2 & -1 & 0 & \cdots & 0 \\ 0 & -1 & 2 & -1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \ddots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & 1 \end{bmatrix} \cdot \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ \vdots \\ v_k \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ \vdots \\ 0 \end{bmatrix}.$$

Then we use **scipy.minimize** on a uniform grid.

```
7 M = 100
_{8}h = 1 / M
9 regularization_lambda = 0.01
11
12 def finite_dimensional_objective(y: np.ndarray) -> float:
      average_as_vector = np.ones(y.shape)
13
      objective = 0.5 * h * np.sum(
14
          np.square(y - average_as_vector)
      ) + 0.5 * regularization_lambda * math.pow(y[0], 2)
16
      return objective
17
18
19
20 def build_constraint_matrix() -> sparse.lil_array:
      constraint_matrix = sparse.lil_array((M, M))
21
      constraint_matrix[-1, -1] = 1
22
      for i in range(1, M - 1):
          constraint_matrix[i, i - 1] = -1
24
          constraint_matrix[i, i] = 2
25
          constraint_matrix[i, i + 1] = -1
26
      return sparse.csr_matrix(constraint_matrix)
27
28
29
30 def build_constraint_vector() -> np.ndarray:
      return np.zeros(M)
31
32
33
34 constraint = opt.LinearConstraint(
      build_constraint_matrix(),
35
      build_constraint_vector(),
36
      build_constraint_vector(),
37
38)
39
40 result = opt.minimize(
      finite_dimensional_objective,
41
      np.ones(M),
42
      method="trust-constr",
43
      constraints=constraint,
44
45)
```

This results in the numerical solution for M = 100 plotted in figure 1, as well as the error for finer and finer grids plotted in figure 2.



Figure 1: Numerical approximation of v = S(u) from scipy.optimize



Figure 2: L_2 grid-norm error for scipy.optimize

Interestingly, we do not see that the error $\rightarrow 0$ as $h \rightarrow 0$ as we would expect.

3 An uncoupled distributed control problem.

The next problem involved a similar problem where the control is *distributed* instead of on the boundary:

Problem 3.1 – 1D uncoupled distributed control problem.

As before, we have a 1-dimensional rod represented by the interval [0, 1], but we fix its temperature to be 0 at both ends and instead control the heat sources in the interior of the rod. Thus, where before the control u was the temperature of the rod on the left endpoint, now $u \in L^2(\Omega)$ defines the heat sources. We still want the temperature of the rod to be as close to 1, pointwise, as possible. This leads the following constrained optimization problem:

minimize
$$J(v, u) := \frac{1}{2} \int_{\Omega} (v(x) - 1)^2 dx + \frac{\lambda}{2} \int_{\Omega} u(x)^2 dx$$
 (4a)

subject to $-\Delta v = u \text{ on } \Omega$, (4b)

$$v(0) = v(1) = 0 \tag{4c}$$

$$\iota \in L^2(\Omega). \tag{4d}$$

 $a \le u \le b$ a.e. in Ω (4e)

Note that the additional constraint on u in (4e) is a typical one and called a *box* constraint.

Here we will not solve the problem analytically, but rather numerically via an implementation of the active set algorithm described in [2].

As a precursor to the algorithm, we start by discretizing the control, u. As above, we partition the domain into subdomains ω_k and approximate u by a piecewise constant function $u_h = \sum_{k=1}^{M} u_k \mathbb{1}_{\omega_k}, u_1, \ldots, u_M \in \mathbb{R}$. We then re-define the solution operator $S : L^2(\Omega) \to L^2(\Omega)$ that maps the control u to the unique solution v to the constraints in (4b)-(4c).

The main difficulty in implementing the algorithm involves finding a discretized version $S_h : \mathbb{R}^M \to L^2(\omega)$ of the solution operator $S : L^2(\Omega) \to L^2(\Omega)$. We do this by computing $v_k := S(\mathbb{1}_{\omega_k})$ for each $k = 1, \ldots, M$, that is, we compute the unique solution v_k to the boundary value problem

$$-\Delta v_k = \mathbb{1}_{\omega_k} \text{ on } \Omega, \tag{5a}$$

$$v_k(0) = v_k(1) = 0.$$
 (4c)

We could very well numerically approximate the solution to this problem using the **ELLIPTIC1d** code used in part (a) of the course, but we can also find an analytic solution without too much difficulty. In 1D, (5a) is equivalent to $-v'' = \mathbb{1}_{\omega_k}$, so we begin by finding a second antiderivative for $\mathbb{1}_{\omega_k}$. Define $f : [0, 1] \to \mathbb{R}$ by

$$f(x) = \int_0^x \mathbb{1}_{\omega_k}(t) dt + C, \quad C \in \mathbb{R}$$
$$= \begin{cases} C & \text{if } x < x_{k-1/2} \\ C + x - x_{k-1/2} & \text{if } x \in \omega_k \\ C + h_k & \text{if } x > x_{k+1/2} \end{cases}$$

so that $f' = \mathbb{1}_{\omega_k}$ by the fundamental theorem of calculus. Then define $g: [0,1] \to \mathbb{R}$ by

$$g(x) = \int_0^x f(t) dt + D, \quad D \in \mathbb{R}$$

=
$$\begin{cases} Cx + D & \text{if } x < x_{k-1/2} \\ Cx + D + x^2/2 + x_{k-1/2}^2/2 - x \cdot x_{k-1/2} & \text{if } x \in \omega_k \\ Cx + D + h_k(x - x_k) & \text{if } x > x_{k+1/2} \end{cases}$$

so that g' = f by the fundamental theorem of calculus. Thus $g'' = \mathbb{1}_{\omega_k}$, so v = -g satisfies (5a). Plugging in boundary conditions yields D = 0, $C = -h_k(1 - x_k)$.

With this formula for v_k , we can define the discretized version of the solution operator $S_h : \mathbb{R}^M \to L^2(\Omega)$ by $S_h(\vec{u}) = \sum_{k=1}^M u_k v_k$. We also need the adjoint $S_h^* : L^2(\Omega) \to \mathbb{R}^n$ of this operator, which, as described on p. 100 in [2], is given by

$$\left(S_h^*(y)\right)_k = \int_\Omega y \cdot v_k \, dx.$$

We will use the quad integration function from scipy.integrate to compute this integral as necessary.

Now, as in [2], the active set algorithm for this problem is as follows:

- **D0** Choose arbitrary initial vectors $\vec{u}_0, \vec{p}_0 \in \mathbb{R}^M$.
- **D1** Given \vec{u}_n and \vec{p}_n , determine the new *finite active sets* A_{n+1}^+ and A_{n+1}^- as well as the finite inactive set I_{n+1} :

$$A_{n+1}^{+} := \left\{ k \in \{1, \dots, M\} : -\frac{p_{n,k}}{\lambda} > b \right\}$$
$$A_{n+1}^{-} := \left\{ k \in \{1, \dots, M\} : -\frac{p_{n,k}}{\lambda} < a \right\}$$
$$I_{n+1} := \{1, \dots, M\} \setminus \left(A_{n+1}^{+} \cup A_{n+1}^{-}\right).$$

D2 Compute $\vec{p}_{n+1} := S_h^* (S_h \vec{u}_n - \mathbb{1}_\Omega)$ and

$$u_{n+1,k} := \begin{cases} a & \text{if } k \in A_{n+1}^{-} \\ -\frac{p_{n+1,k}}{\lambda} & \text{if } k \in I_{n+1} \\ b_k & \text{if } k \in A_{n+1}^{+}. \end{cases}$$

Proceed to **D1** for the next iteration step.

The iteration is halted once $A_{n+1}^+ = A_n^+$ and $A_{n+1}^- = A_n^-$. This algorithm is implemented as follows:

```
43 def analytic_v_k(x: float, k: int = 0) -> float:
      C = -cell_widths[k] * (1 - cell_centers[k])
44
      if x <= cell_boundaries[k]:</pre>
45
           return -C * x
46
      elif cell_boundaries[k] <= x <= cell_boundaries[k + 1]:</pre>
47
          return -(
48
               0.5 * math.pow(x, 2)
49
               - x * cell_boundaries[k]
               + 0.5 * math.pow(cell_boundaries[k], 2)
               + C * x
      return -(cell_widths[k] * (x - cell_centers[k]) + C * x)
57 def integral_of_v_k(k: int) -> float:
      return quad(analytic_v_k, 0, 1, args=(k))[0]
58
61 integrals_of_v_k = np.array([integral_of_v_k(i) for i in range(M)])
62
63 initial_u = np.ones(M)
64 initial_p = np.ones(M)
_{65} p, \upsilon = initial_p, initial_\upsilon
66
67 (
      previous positive active set,
68
      previous_negative_active_set,
69
      previous_inactive_set,
70
      positive_active_set,
71
      negative_active_set,
72
      inactive_set,
73
74) = (set(), set(), set(), set(), set(), set())
_{75} n = 0
76 while n < 1000:
      n += 1
      print("iteration " + str(n))
78
      (
80
          previous_positive_active_set,
81
          previous_negative_active_set,
82
          previous_inactive_set,
83
      ) = (positive_active_set, negative_active_set, inactive_set)
84
      positive_active_set, negative_active_set, inactive_set = set(), set()
85
      , set()
86
      for k in range(M):
87
          if -p[k] / regularization_lambda > b[k]:
88
               positive_active_set.add(k)
89
          elif -p[k] / regularization_lambda < a[k]:</pre>
90
               negative_active_set.add(k)
91
```

```
else:
92
               inactive_set.add(k)
93
94
      p = np.zeros(M)
95
       for k in range(M):
96
           for j in range(M):
               p[k] += u[j] * integrals_of_v_k[j]
98
           p[k] -= integrals_of_v_k[k]
99
100
      for k in positive_active_set:
           u[k] = a[k]
102
       for k in negative_active_set:
           v[k] = b[k]
104
      for k in inactive_set:
105
           u[k] = -p[k] / regularization_lambda
106
107
      if (
108
           positive_active_set == previous_positive_active_set
109
           and negative_active_set == previous_negative_active_set
110
      ):
           print("Found u: " + str(u))
           break
113
```

For $a \equiv 0$ and $b \equiv 10$, this produced the result of $u_k = 10$ for all $k = 1, \ldots, M$, and thus v(x) = -5x(x-1).

4 Future work.

For C3, I would like to work on a coupled problem. For instance, if $\Omega_1 = [0, 1/2]$, $\Omega_2 = [1/2, 1]$, and $v_1 \in L^2(\Omega_1)$, $v_2 \in L^2(\Omega_2)$, we can try to control $u \in \mathbb{R}$ such that v_2 is as close to 1, pointwise, as possible, and such that

 $-\Delta v_i = 0, \quad v_1(0) = u, \quad v_1(1/2) = v_2(1/2), \quad v_1'(1/2) = v_2'(1/2), \quad v_2(1) = 0.$

Other future directions which I may not have time for would (1) include implementing my own algorithm for the boundary control problems or (2) solving a 2D distributed control problem.

References

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- [2] P. Philip, "Optimal control of partial differential equations," 2013. [Online]. Available: https://www.math.lmu.de/~philip/publications/lectureNotes/philipPeter_ OptimalControlOfPDE.pdf